

# Chap 7. Regularization

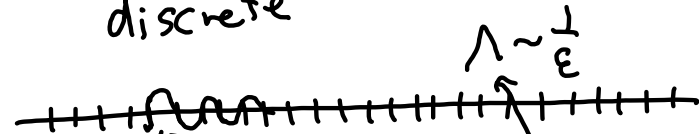
continuum



$$\sum_{\lambda \rightarrow \text{wave length}} \rightarrow \int d^3 \vec{k}$$

→

discrete



$$\int_{|\vec{k}| = \Lambda \rightarrow \infty}$$

cutoff

---

dimensional regularization

# LSZ reduction formula

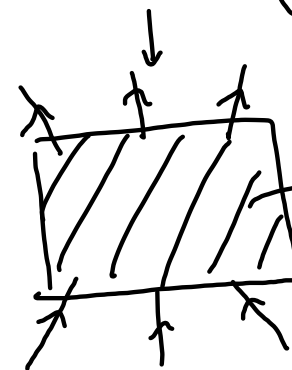
Correlation function

$$\langle \Omega | T \{ \phi(x_1) \dots \phi(x_n) \} | \Omega \rangle$$



Scattering matrix

$$\langle \vec{p}_1, \dots | \vec{k}_1, \dots \rangle \equiv \langle \vec{p}_i | S | \vec{k}_j \rangle$$



on-shell  
 $p_i^2 = E_i^2 - \vec{p}_i^2 = m_i^2$

F.D.  
 (connected amputated)

$$\int \dots \int \langle \Omega | T \{ \phi(x_1) \dots \phi(x_n) \phi(y_1) \dots \phi(y_m) \} | \Omega \rangle \prod_{i=1}^n e^{i p_i \cdot x_i} \prod_{j=1}^m e^{-i k_j \cdot y_j} d^4 x_i d^4 y_j$$

$$\sim \prod_i \left( \frac{i \sqrt{Z}}{p_i^2 - m_i^2 + i \epsilon} \right) \prod_j \left( \frac{i \sqrt{Z}}{k_j^2 - m_j^2 + i \epsilon} \right) \langle \vec{p}_i | S | \vec{k}_j \rangle$$

on-shell.

$$p_i^2 = m_i^2$$

$$k_j^2 = m_j^2$$

Any state  $\rightarrow \mathbb{1} = \sum_n |n\rangle \langle n|$

Q.M.  $\rightarrow \mathbb{1}_{1\text{-particle}} = \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2E_p} |\vec{p}\rangle \langle \vec{p}|$

QFT  $\mathbb{1} = |\Omega\rangle \langle \Omega| + \sum_\lambda \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2E_p(\omega)} |\lambda_{\vec{p}}\rangle \langle \lambda_{\vec{p}}|$

$\vec{p}$  = total momentum of multi-particle states

(ex) 1-particle:  $p = (p^0, \vec{p})$   $p^0 = \sqrt{\vec{p}^2 + m^2} = E_{\vec{p}}$

2- " :  $p = p_1 + p_2 = (\underbrace{p_1^0 + p_2^0}_{E_p}, \underbrace{\vec{p}_1 + \vec{p}_2}_{\vec{p}})$   $E_p = \sqrt{\vec{p}_1^2 + m^2} + \sqrt{\vec{p}_2^2 + m^2}$   
 $= \underbrace{E_{\vec{p}}(\lambda)}_{\leftarrow \vec{p}_1 - \vec{p}_2} = \sqrt{(\vec{p}_1 + \vec{p}_2)^2 + m^2}$

$$\begin{aligned} \langle \Omega | \phi(x) \phi(y) | \Omega \rangle &= \langle \Omega | \phi(x) | \Omega \rangle \langle \Omega | \phi(y) | \Omega \rangle + \sum_\lambda \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2E_p(\omega)} \langle \Omega | \phi(x) | \lambda_{\vec{p}} \rangle \langle \lambda_{\vec{p}} | \phi(y) | \Omega \rangle \\ &= \langle \Omega | \phi(0) | \Omega \rangle \langle \Omega | \phi(0) | \Omega \rangle e^{+i p \cdot y} e^{-i E_p x^0 + i \vec{p} \cdot \vec{x}} \\ &= \langle \Omega | \phi(0) | \Omega \rangle \langle \Omega | \phi(0) | \lambda_{\vec{p}} \rangle e^{-i p \cdot x} \Big|_{p^0 = E_p(\lambda)} \end{aligned}$$

$$= \sum_{\lambda} \int \frac{d^3 p}{(2\pi)^3} \frac{1}{2E_{\mathbf{p}}(\lambda)} e^{-i\mathbf{p} \cdot (\mathbf{x}-\mathbf{y})} \left| \langle \Omega | \phi(\omega) | \lambda_0 \rangle \right|^2$$

$\uparrow$   $D_F(x-y, m_{\lambda}^2) = \int dM^2 \delta(M^2 - m_{\lambda}^2) D_F(x-y, M^2)$

$p^0 = E_{\mathbf{p}}(\lambda) = \sqrt{\mathbf{p}^2 + m_{\lambda}^2}$

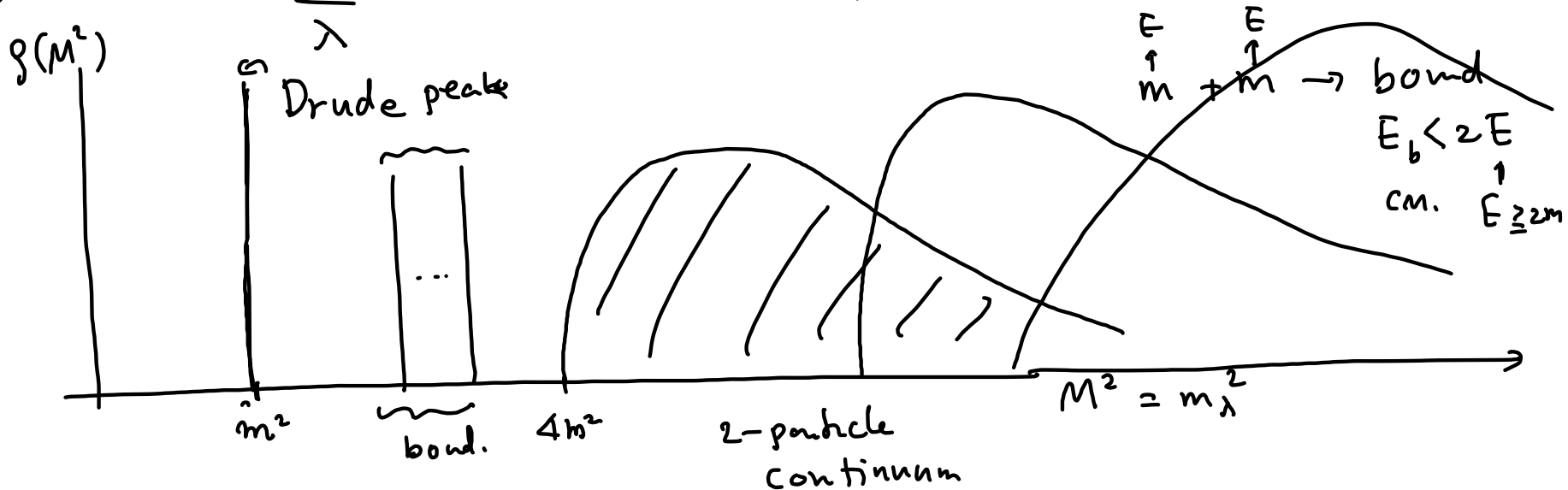
$$= \sum_{\lambda} \int \frac{d^4 p}{(2\pi)^4} \frac{i}{p^2 - m_{\lambda}^2 + i\epsilon} e^{-i\mathbf{p} \cdot (\mathbf{x}-\mathbf{y})} \left| \langle \Omega | \phi(\omega) | \lambda_0 \rangle \right|^2$$

$\uparrow$  off-shell

spectral density function

$$= \int \frac{dM^2}{2\pi} \rho(M^2) D_F(x-y, M^2)$$

$$\rho(M^2) \equiv \sum_{\lambda} (2\pi) \delta(M^2 - m_{\lambda}^2) \left| \langle \Omega | \phi(\omega) | \lambda_0 \rangle \right|^2$$



$$\int d^4x e^{ip \cdot x} \langle \Omega | \phi(x) \phi(0) | \Omega \rangle = \int \frac{dM^2}{2\pi} \rho(M^2) \int D_F(x, M^2) d^4x e^{ip \cdot x}$$

$$= \frac{iZ}{p^2 - m^2 + i\epsilon} + \int_{4m^2}^{\infty} \frac{dM^2}{2\pi} \rho(M^2) \frac{i}{p^2 - M^2 + i\epsilon}$$

$2\pi \delta(M^2 - m^2) Z \approx |\langle \Omega | \phi(0) | \lambda_0 \rangle|^2$

---

Electron self-energy

$$\langle \Omega | T \{ \psi(x) \bar{\psi}(y) \} | \Omega \rangle_{\text{QED}} = \text{diagram 1} + \text{diagram 2} + \dots$$

momentum space

$$= \text{diagram 1} + \text{diagram 2}$$

$$= \frac{i(\not{p} + m_0)}{p^2 - m_0^2} + \frac{i(\not{p} + m_0)}{p^2 - m_0^2} (-i\Sigma_2(p)) \frac{i(\not{p} + m_0)}{p^2 - m_0^2}$$

$$-i \Sigma_2(p) \equiv (-ie)^2 \int \frac{d^4 k}{(2\pi)^4} \gamma^\mu \frac{i(\not{k} + m_0)}{\underbrace{k^2 - m_0^2 + i\epsilon}_B} \gamma^\mu \frac{-i}{\underbrace{(p-k)^2 - \mu^2 + i\epsilon}_A}$$

Trick "Feynman parametrization"

$$\begin{aligned} \frac{1}{AB} &= \int_0^1 dx \frac{1}{[Ax + B(1-x)]^2} = \int_0^1 dx \frac{1}{((A-B)x + B)^2} \\ &= - \frac{1}{(A-B)x + B} \cdot \frac{1}{A-B} \Big|_0^1 \\ &= - \frac{1}{A-B} \cdot \left( \frac{1}{A} - \frac{1}{B} \right) = \frac{1}{AB} \quad \checkmark \end{aligned}$$

---


$$\Rightarrow \frac{1}{\prod_{i=1}^n A_i} = \int_0^1 \dots \int_0^1 \frac{1}{\left( \sum_{i=1}^n A_i x_i \right)^{n+1}} \delta\left(\sum_i x_i = 1\right) dx_1 \dots dx_n$$

$$Ax + B(1-x) = \left( \underset{\uparrow}{(p-k)^2} - \mu^2 + i\varepsilon \right) x + \left( \underset{\uparrow}{k^2} - m_0^2 + i\varepsilon \right) (1-x)$$

$$= \underline{k^2} + p^2 x - \underline{2p \cdot k x} - \mu^2 x - m_0^2 (1-x) + i\varepsilon$$

$$= \underbrace{\left( \underbrace{k - x p}_{\ell} \right)^2}_{d^4 k \rightarrow d^4 \ell} - \underbrace{x^2 p^2 + p^2 x - \mu^2 x - m_0^2 (1-x) + i\varepsilon}_{\Delta^2} + p^2 x (1-x)$$

$$= \ell^2 - \Delta^2 + i\varepsilon$$

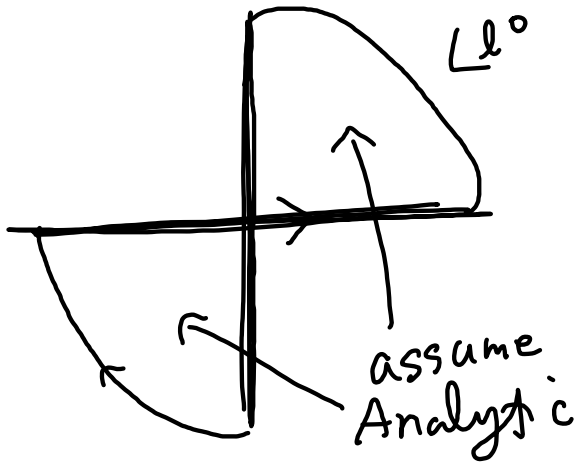
$$-i \Sigma_2(p) = -e^2 \int_0^1 dx \int \frac{d^4 \ell}{(2\pi)^4} \frac{\gamma^\mu \overbrace{(x \cancel{p} + \cancel{\ell} + m_0)}^k \gamma_\mu}{\ell^2 - \Delta^2 + i\varepsilon} \rightarrow \frac{p_\nu \overbrace{\gamma^\mu \gamma^\nu \gamma_\mu}^{-2\gamma^\nu}}{-2x \cancel{p}} + m_0 \overbrace{\gamma^\mu \gamma_\mu}^{4\mathbb{1}}$$

$$\int \frac{d^4 \ell}{\ell^2 - \Delta^2} \cancel{\ell}^{\leftarrow \text{odd } (\ell \rightarrow -\ell)} = 0$$

$$= -e^2 \int_0^1 dx \left( -2x \cancel{p} + 4m_0 \mathbb{1} \right) \int \frac{d^4 \ell}{(2\pi)^4} \frac{1}{\ell^2 - \Delta^2 + i\varepsilon}$$



$$l^2 = l_0^2 - (l_1^2 + l_2^2 + l_3^2)$$



Wick rotation:

$$\int_{-\infty}^{\infty} dl_0 = \int_{-i\infty}^{i\infty} dl_0 = i \int_{-\infty}^{\infty} dl_E$$

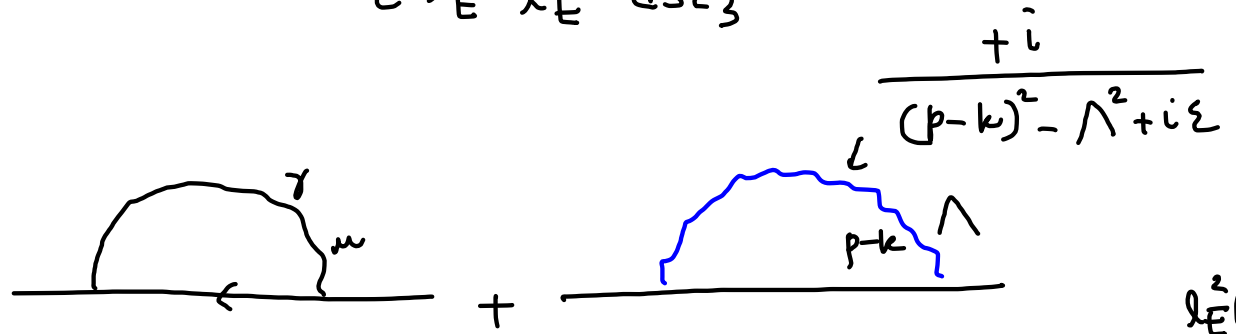
$l_0 = i l_4$

$$l^2 = -(l_1^2 + \dots + l_4^2) \quad l_E = (l_1, l_2, l_3, l_4)$$

$$\int d^4 l \frac{1}{(l^2 - \Delta^2)^2} = -i \int d^4 l_E \frac{1}{(l_E^2 + \Delta^2)^2} = -i \int d\Omega_3 \int_0^{\infty} \frac{l_E^3 dl_E}{(l_E^2 + \Delta^2)^2}$$

## Regularization

### ① Pauli-Villars



가상의 heavy photon

$$\frac{i}{(4\pi)^2} \log \frac{\Delta_\mu^2}{\Delta_\Lambda^2} = \frac{i}{(4\pi)^2} \int_0^\infty dl_E^2 \left[ \frac{l_E^2}{(l_E^2 + \Delta_\mu^2)^2} - (m \rightarrow \Lambda) \right]$$

$$\int d^4 l \left( \frac{1}{(l^2 - \Delta_\mu^2)^2} - \frac{1}{(l^2 - \Delta_\Lambda^2)^2} \right)$$

$\int_0^\infty \frac{l_E^3 dl_E}{(l_E^2 + \Delta_\mu^2)^2} - \int_0^\infty \frac{l_E^3 dl_E}{(l_E^2 + \Delta_\Lambda^2)^2}$



$$-i \Sigma_2(p) = \frac{e^2}{8\pi^2} \int_0^1 dx (2m_0 \not{1} - x \not{p}) \log \left( \frac{x \Lambda^2}{(1-x)m_0^2 + x\mu^2 - x(1-x)p^2} \right)$$

$\uparrow \quad \uparrow \quad \Delta^2$   
 $\not{p} \quad \mu$

$\frac{\partial \Sigma_2}{\partial \not{p}}$

$$-i \Sigma_{\text{all}}(p) = \Sigma(-i \Sigma(p))$$

$$-i \Sigma = \text{[Diagram: wavy line with a blob on top]} + \text{[Diagram: wavy line with a blob on bottom]} + \text{[Diagram: wavy line with a blob on top and bottom]} + \dots$$

$$= \text{[Diagram: circle with '1PI' inside]} \text{---}$$

$$\text{[Diagram: chain of circles with '1PI' inside, connected by arrows]} = \text{[Diagram: chain of circles with '1PI' inside, connected by arrows]} + \text{[Diagram: chain of circles with '1PI' inside, connected by arrows]} + \dots$$

$$= \frac{i(\not{p} + m_0)}{p^2 - m_0^2} \left\{ 1 + \underbrace{(-i \Sigma)}_{\text{wavy}} \frac{i(\not{p} + m_0)}{p^2 - m_0^2} + \left( (-i \Sigma) \frac{i(\not{p} + m_0)}{p^2 - m_0^2} \right)^2 + \dots \right\} = \frac{i}{\not{p} - m_0} \cdot \frac{1}{1 - \frac{\Sigma}{\not{p} - m_0}}$$

$$= \frac{i}{\not{p} - m_0 - \sum(\not{p})}$$

$\uparrow$   
 $4 \times 4$

→

$$\not{p} - m_0 - \sum(\not{p}) \Big|_{\not{p}=m} = 0$$

(CF)

$$\leftarrow = \frac{i}{\not{p} - m_0} = \frac{i(\not{p} + m_0)}{p^2 - m_0^2}$$

$\nwarrow$   
 $\not{p} \sim m_0$

① mass renormalization.

$$\underbrace{\delta m}_{+\infty} = \sum(\not{p}=m) \leftarrow \underbrace{m - m_0}_{\substack{\downarrow 0.5 \text{ MeV} \\ \delta m}} - \underbrace{\sum(\not{p}=m)}_{= c e^2 \ln \Lambda^2} = 0$$

$$- c e^2 \ln \Lambda^2 + \dots$$

② wave function renorm.

$$\sum(\not{p}) \approx \sum(m) + (\not{p} - m) \sum'(m) + \dots$$

$$\frac{i}{\not{p} - m_0 - \underbrace{\sum(m)}_{-m} - (\not{p} - m) \sum'(m)}$$

$$= \frac{i}{(\not{p} - m) \underbrace{(1 - \sum'(m))}_{\frac{1}{Z_2}})}$$

$$= \frac{i Z_2}{\not{p} - m}$$

$$\frac{1}{Z_2} = \frac{1}{1 + \delta Z_2} = 1 - \delta Z_2$$

1-loop.

$$\alpha = \frac{e^2}{4\pi}$$

$$\delta m = \frac{\alpha}{2\pi} m_0 \int_0^1 dx (2-x) \ln \left[ \frac{x \Lambda^2}{(1-x)^2 m_0^2 + x \mu^2} \right].$$

$$\overline{Z}_2 = 1 + \delta Z_2$$

$$\delta Z_2 = \Sigma'(m) = \left. \frac{\partial \Sigma(\not{p})}{\partial \not{p}} \right|_{\not{p}=m} = \frac{\alpha}{2\pi} \int_0^1 dx [ \dots ]$$

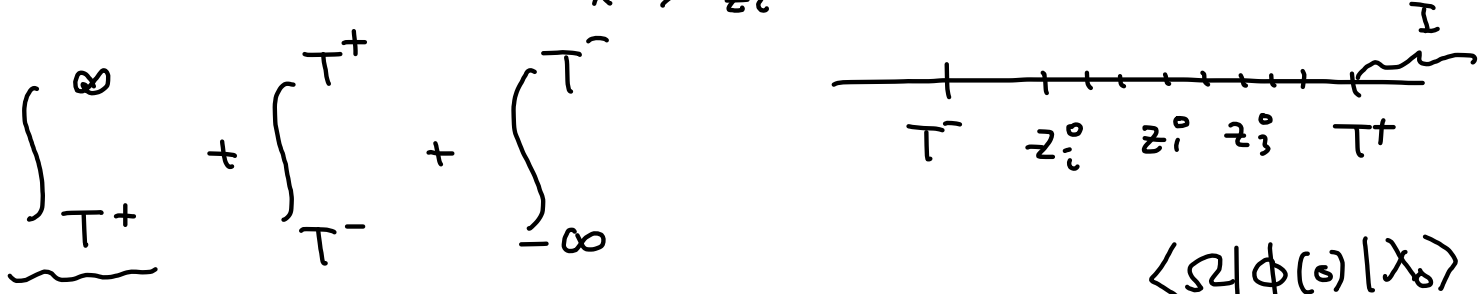
$$\int d^4x \langle \Omega | T \{ \phi(x) \phi(\omega) \} | \Omega \rangle \underset{p^2 \rightarrow m^2}{=} \frac{i Z}{p^2 - m^2 + i\epsilon}$$

# LSZ reduction

$$\int d^4x e^{i p \cdot x} \langle \Omega | T \{ \underbrace{\phi(x)} \underbrace{\phi(z_1) \phi(z_2) \dots} \} | \Omega \rangle$$

$$\int d^3\vec{x} \int_{-\infty}^{\infty} dx^0 \quad \underline{=} \quad \langle \Omega | \phi(x) \overset{\downarrow}{T} \{ \phi(z_1) \dots \} | \Omega \rangle$$

$x^0 > z_i^0$



$$= \int d^3\vec{x} \int_{T^+}^{\infty} dx^0 \quad e^{i p^0 x^0} e^{-i \vec{p} \cdot \vec{x}} \sum_{\lambda} \int \frac{d^3\vec{q}}{(2\pi)^3} \frac{1}{2E_{\vec{q}}(\lambda)} \underbrace{\langle \Omega | \phi(x) | \lambda_{\vec{q}} \rangle}_{\delta^0 = E_{\vec{q}}(\lambda)}$$

$$\times \langle \lambda_{\vec{q}} | T \{ \phi(z_1) \dots \} | \Omega \rangle$$

$$\underbrace{e^{-i \vec{q} \cdot \vec{x}} e^{-\epsilon x^0}}_{\epsilon \rightarrow 0^+} \int d^3\vec{x} = \int_{T^+}^{\infty} dx^0 e^{i x^0 (p^0 - \delta_{\vec{q}}^0 + i\epsilon)} \underbrace{E_{\vec{q}}(\lambda)}_{\delta^0 = E_{\vec{q}}(\lambda)}$$

$$(2\pi)^3 \delta^{(3)}(\vec{p} - \vec{q})$$

$$\Rightarrow \sum_{\lambda} \frac{1}{2E_{\vec{p}(\lambda)}} \frac{i}{p^0 - E_{\vec{p}(\lambda)} + i\epsilon} e^{i(p^0 - E_{\vec{p}} + i\epsilon)T^+} \underbrace{\langle \Omega | \phi(z) | \lambda_0 \rangle \langle \lambda_{\vec{p}} | T \{ \phi(z) \dots \} | \Omega \rangle}_{\sqrt{Z}}$$

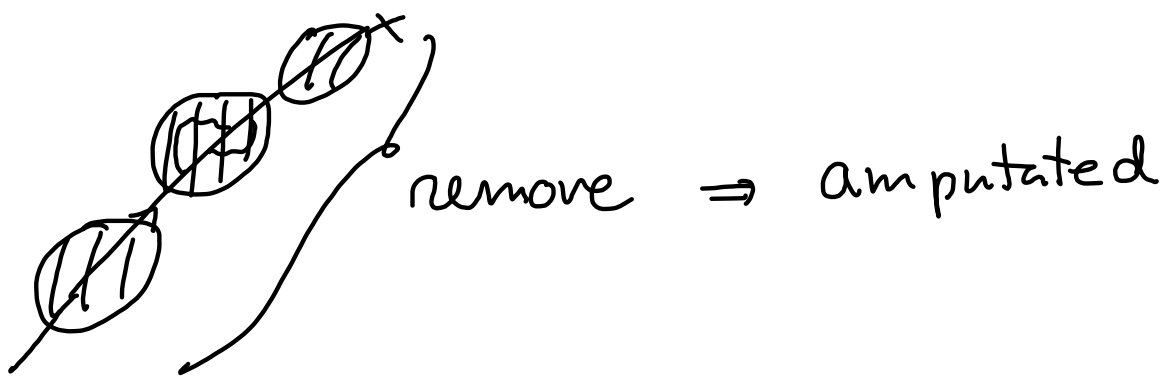
$$\frac{i}{p^2 - m_{\lambda}^2 + i\epsilon} \Big|_{p^2 = m_{\lambda}^2} = \frac{i}{p^0 - (E_{\vec{p}} + i\epsilon)(p^0 + E_{\vec{p}})} \approx \frac{i}{2E_{\vec{p}}(p^0 - E_{\vec{p}})}$$

$$= \sum_{\lambda} \frac{i\sqrt{Z}}{p^2 - m_{\lambda}^2 + i\epsilon} \langle \lambda_{\vec{p}} | T \{ \phi(z_1) \dots \} | \Omega \rangle$$

$$= \frac{i\sqrt{Z}}{p^2 - m^2 + i\epsilon} \langle \vec{p} | T \{ \phi(z_1) \dots \} | \Omega \rangle + \text{continuum}$$

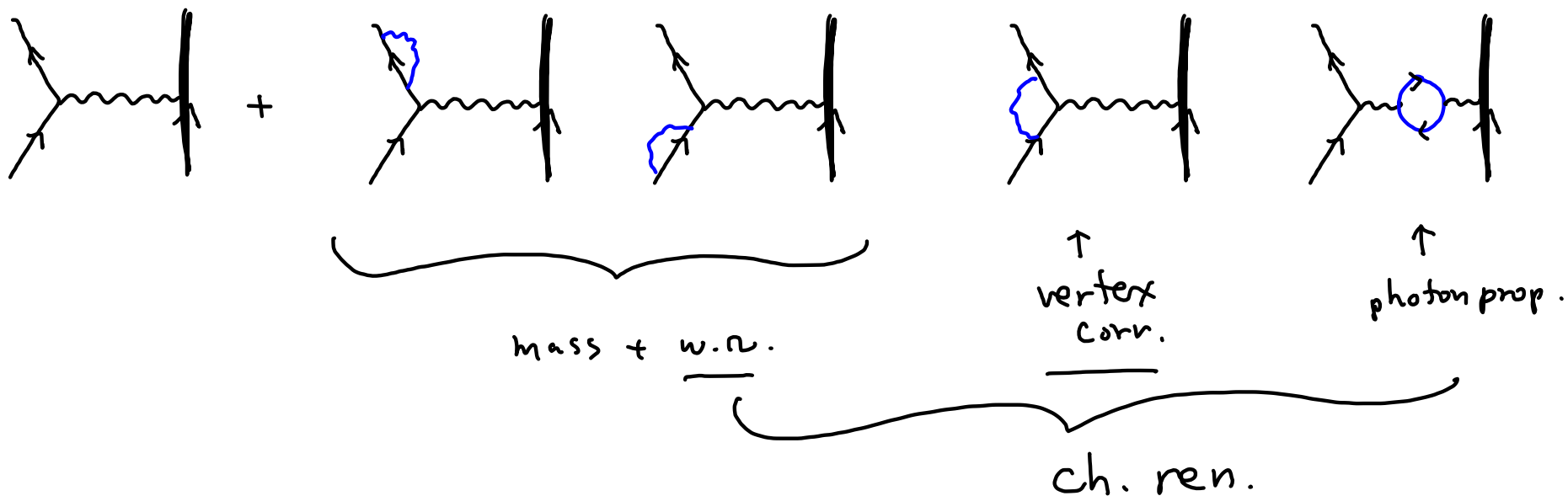
$$\dots \frac{\prod_i \frac{i\sqrt{Z}}{i p_i^2 - m^2 + i\epsilon}}{\prod_j \frac{i\sqrt{Z}}{k_j^2 - m^2 + i\epsilon}} \langle \vec{p}_1, \vec{p}_2, \dots | T \{ \dots \} | k_1, \dots, k_m \rangle_{in} = \left( \langle \Omega | T \{ \phi(x_1) \dots \phi(x_n) \} | \Omega \rangle \right) e^{i p \cdot x} e^{-i k \cdot y}$$

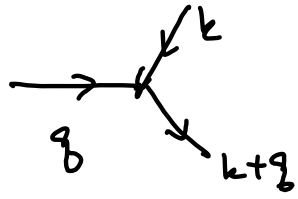
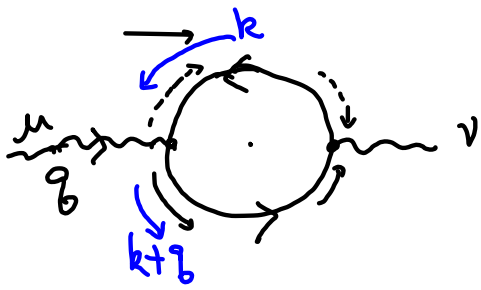
S-matrix



$$QED : e_0, m_0 \rightarrow e, m$$

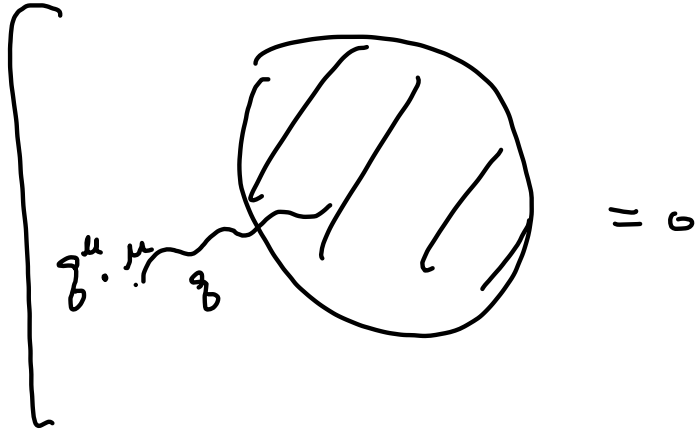
③ charge renorm.





$$= (-ie)^2 (-1) \int \frac{d^4 k}{(2\pi)^4} \text{tr} \left[ \gamma^\mu \frac{i}{\not{k}-m} \gamma^\nu \frac{i}{\not{k}+\not{q}-m} \right]$$

$$\equiv i \Pi_2^{\mu\nu}(q) = i (q^2 g^{\mu\nu} - q^\mu q^\nu) \underline{\underline{\Pi_2(q)}}$$



Ward-Takahashi Identity

$$q_\mu \Pi_2^{\mu\nu} = q_\nu \Pi_2^{\mu\nu} = 0$$


---


$$q^2 q^\nu + \alpha q^2 q^\nu = 0$$

$$\alpha = -1$$

$$\text{loop} + \text{loop with cut} + \text{loop with blob} + \dots = \text{loop with 1PI blob} = i \Pi^{\mu\nu}$$

$$= \underline{i (q^2 g^{\mu\nu} - q^\mu q^\nu) \Pi(q)}$$

$$\text{loop} + \boxed{\text{1PI}} + \text{loop with 1PI} + \dots = \frac{-i g_{\mu\nu}}{q^2} + \frac{-i g_{\mu\rho}}{q^2} (i \Pi^{\rho\sigma}) \frac{-i g_{\sigma\nu}}{q^2} + \dots$$

$$= \frac{-i}{g^2(1-\Pi(g^2))} \left( g_{\mu\nu} - \frac{g_{\mu}g_{\nu}}{g^2} \right) + \frac{-i}{g^2} \left( \frac{g_{\mu}g_{\nu}}{g^2} \right)$$

$\therefore g_{\mu}(\quad) \sim g^2 = 0$

$$= \frac{-i g_{\mu\nu}}{g^2(1-\Pi(g^2))} = \frac{-i g_{\mu\nu}}{g^2(1-\Pi(0) + \dots)} = \frac{-i g_{\mu\nu} Z_3}{g^2}$$

$$Z_3 = \frac{1}{1-\Pi(0)}$$

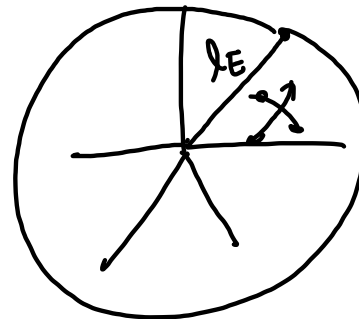


# Dimensional Reg.

$$4 \rightarrow d$$

$$\int \frac{d^4 l_E}{(l_E^2 + \Delta^2)^2} \sim \int \frac{l_E^4 dl_E}{l_E^4 l_E} \sim \text{logarithmic Div.}$$

$$\int \frac{d^d l_E}{(2\pi)^d} = \int \frac{d\Omega_d}{(2\pi)^d} \int_0^\infty l_E^{d-1} dl_E \dots$$



$$(\sqrt{\pi})^d = \left( \int_{-\infty}^{\infty} dx e^{-x^2} \right)^d = \int_{-\infty}^{\infty} \dots \int_{-\infty}^{\infty} dx_1 \dots dx_d e^{-\underbrace{(x_1^2 + x_2^2 + \dots + x_d^2)}_{x^2}} = \int d\Omega_d \int_0^\infty x^{d-1} dx e^{-x^2}$$

$$= \int d\Omega_d \int_0^\infty dt \frac{t^{\frac{d-1}{2} - \frac{1}{2}}}{2} e^{-t}$$

$\underbrace{\hspace{10em}}_{\Gamma(\frac{d}{2})}$

$$x^2 \equiv t \rightarrow 2x dx = dt$$

$$dx = \frac{dt}{2\sqrt{t}}$$

$$\therefore \int d\Omega_d = \frac{2\pi^{d/2}}{\Gamma(\frac{d}{2})}$$

$$d=3 \quad 4\pi = \frac{2\pi^{3/2}}{\Gamma(\frac{3}{2})} = \frac{2\pi^{3/2}}{\frac{1}{2}\Gamma(\frac{1}{2})} = \frac{2\pi^{3/2}}{\frac{1}{2} + 1 \sqrt{\pi}}$$

$$d=4; \quad \frac{2\pi^2}{\Gamma(2)} = 2\pi^2 //$$

$$\int \frac{1}{(2\pi)^d} \frac{d^d l}{(l^2 + \Delta)^2} = \int d\Omega_d \int_0^\infty \frac{l^{d-2} \sqrt{l dl} \rightarrow \frac{1}{2} d l^2}{(l^2 + \Delta)^2}$$

$$d l^2 = 2 l dl$$

$$\Delta^{\frac{d}{2}-2} \Gamma(2-\frac{d}{2}) \pi^{d/2} \frac{1}{(2\pi)^d} = \int d\Omega_d \frac{1}{2} \int_0^\infty \frac{l^{\frac{d-2}{2}} d l^2}{(l^2 + \Delta)^2}$$

$\frac{1}{2} \frac{\pi^{d/2}}{\Gamma(\frac{d}{2})} \times \Delta^{\frac{d-2}{2}-1} \frac{\Gamma(\frac{d}{2}) \Gamma(-\frac{d}{2}+2)}{\Gamma(2)}$

$\Delta^{\frac{d-2}{2}-1} \int_0^1 dx (1-x)^{\frac{d-2}{2}-1} x^{-\frac{d}{2}+2-1} = \int_0^1 dx (1-x)^{\frac{d-2}{2}-1} x^{-\frac{d-2}{2}}$

$\beta = \frac{d-2}{2}, \alpha = \frac{d-2}{2}$

$$\frac{\Delta}{l^2 + \Delta} \equiv x \rightarrow l^2 + \Delta = \frac{\Delta}{x}$$

$$l^2 = \frac{\Delta(1-x)}{x} \quad dl^2 = -\frac{\Delta}{x^2} dx$$

$$\int_0^1 \frac{\left(\Delta \frac{(1-x)}{x}\right)^{\frac{d-2}{2}} \left(-\frac{\Delta}{x^2}\right) dx}{\left(\frac{\Delta}{x}\right)^2}$$

$$\int_0^1 dx x^{\alpha-1} (1-x)^{\beta-1} = \frac{\Gamma(\alpha) \Gamma(\beta)}{\Gamma(\alpha+\beta)}$$

$$\Delta^{\frac{d}{2}-2} \Gamma\left(2-\frac{d}{2}\right) \pi^{d/2} \frac{1}{(2\pi)^d}$$

$$d \rightarrow 4 \quad \Gamma(0) = \infty$$

$$4-d \equiv \epsilon$$

$$\Gamma\left(2-\frac{d}{2}\right) = \Gamma\left(\frac{\epsilon}{2}\right) \sim \frac{2}{\epsilon} - \gamma + O(\epsilon)$$

$$\pi^{d/2} = \pi^{2-\frac{\epsilon}{2}}$$

$\uparrow$   
Euler-Mascheroni  
0.57721...

$$\Delta^{-\frac{\epsilon}{2}} \approx e^{\log(\Delta^{-\frac{\epsilon}{2}})} = e^{-\frac{\epsilon}{2} \log \Delta} \approx 1 - \frac{\epsilon}{2} \log \Delta + \dots$$

$$= \frac{1}{2^4 \pi^2} \left( \frac{2}{\epsilon} - \gamma + \dots \right) \left( 1 - \frac{\epsilon}{2} \log \Delta + \dots \right)$$

$$\frac{2}{\epsilon} - \gamma - \underline{\underline{\log \Delta}} + \dots$$

$$\int \frac{d^d l_E}{(2\pi)^d} \frac{1}{(l_E^2 + \Delta)^n} = \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(n - \frac{d}{2})}{\Gamma(n)} \left(\frac{1}{\Delta}\right)^{n - \frac{d}{2}}$$

$$\int \frac{d^d l_E}{(2\pi)^d} \frac{l_E^2}{(l_E^2 + \Delta)^n} = \frac{d/2}{(4\pi)^{d/2}} \frac{\Gamma(n - \frac{d}{2} - 1)}{\Gamma(n)} \left(\frac{1}{\Delta}\right)^{n - \frac{d}{2} - 1}$$

$$\gamma^\mu \gamma^\nu \gamma_\mu = -2 \gamma^\nu \quad d=4$$

$$\{\gamma^\mu, \gamma^\nu\} = 2g^{\mu\nu} \mathbb{1} = \frac{- (d-2) \gamma^\nu}{\quad}$$

$$g^{\mu\nu} g_{\mu\nu} = \delta^\mu_\mu = d$$

$$\Rightarrow \Pi_2 = \int \frac{d^d l_E}{(2\pi)^d} \frac{(1 - \frac{d}{2}) g^{\mu\nu} l_E^2}{(l_E^2 + \Delta)^2}$$

$$\Pi_2^{(g^2)} = \frac{-8 p^2}{(4\pi)^{d/2}} \int_0^1 dx x(1-x)$$

$$\Gamma(2 - \frac{d}{2}) \left(\frac{1}{\Delta}\right)^{2 - \frac{d}{2}}$$

$$\left( \frac{2}{\epsilon} - \log \mu^2 - \gamma + \log(4\pi) \dots \right)$$